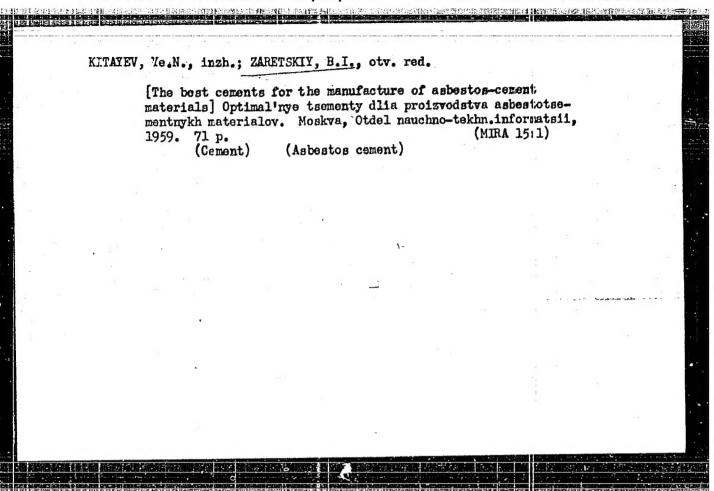
ZARETSKI, B. I.
Tractors. Moskva, Gss. nauchno-tekhn. izd-vo mashinostroit.lit-ry, 1950. 135 p. (50-55251)

TL233.23



KITAYEV, Ye.N., inzh.; GONCHARSKAYA, R.E.; ZAHETSKIY, B.I., ctv. red.; ERLIKH, I.A., red.

[Asbestos cement materials obtained from sand cements by autocleve treatment, and their chemical resistance to corrosive solutions [Khimcheskaia stotkost' v agressivnykh rastvorakh asbestotsementnykh materialov, poluchaemykh iz poschanistykh tsementov a primeneniem avtoklavnoi obrabotki. Moskva, Otdel nauchno-tekhn. informatsii, 1960. 24 p.

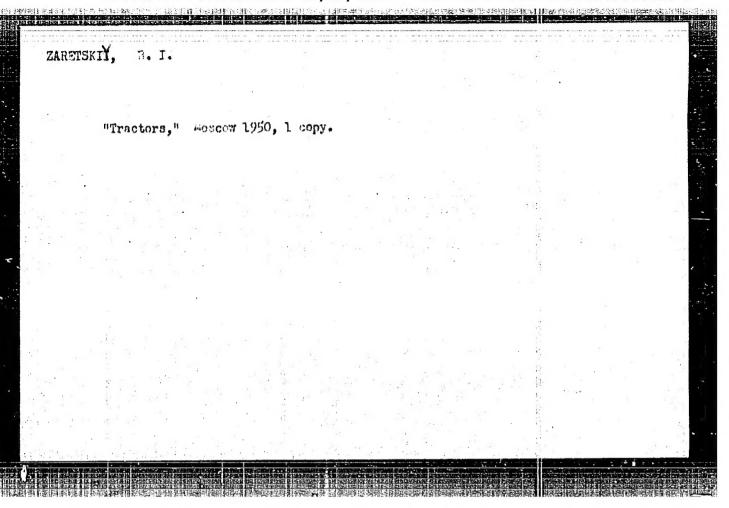
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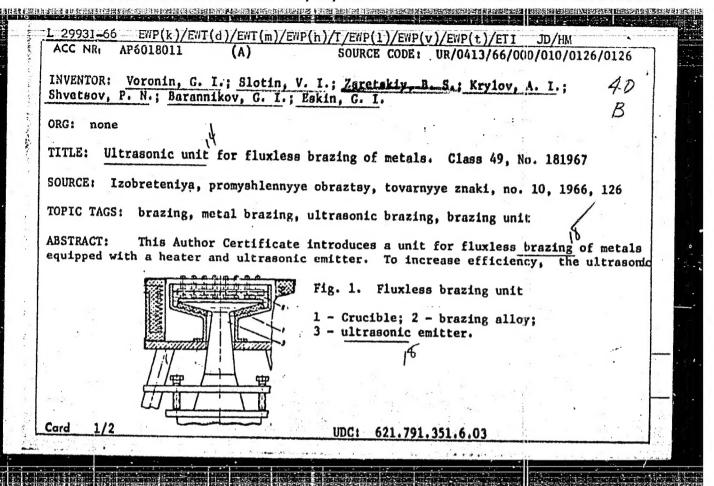
(Asbestos cement)

KAREL'SKIKH, D.K., prof.; APASHEV, M.D., kand.tekhn.nauk; BAHSKIY, I.B., kand.tekhn.nauk; ZAYCHIK, G.I., doktor tekhn.nauk, retsensent; ANOKHIN, V.I., kand.tekhn.nauk, retsensent; ZARETSKIY, B.I.,

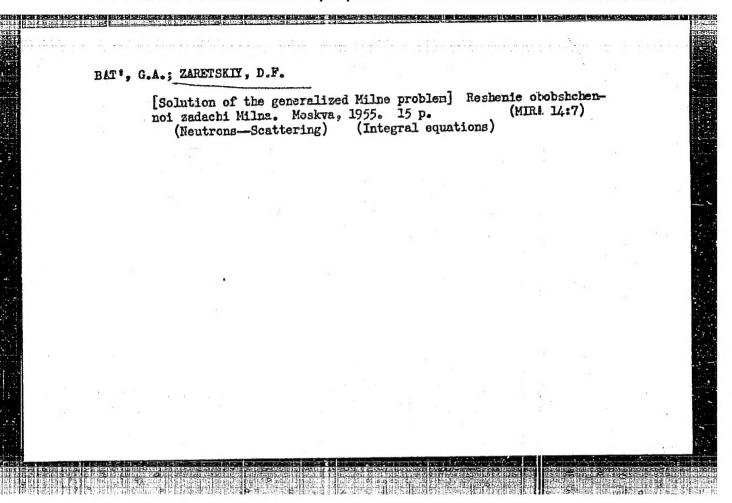
ingh., red.: POPOVA, S.M., tekhn.red.

[Theory, design, and engineering analysis of tractors] Teoriia, konstruktsiia i raschet traktorov. Moskva, Gos.nauchno-tekhn. izd-vo mashinostroit. lit-ry. [Pt.3. Theory and analysis of tractor chassis] Teoriia i raschet shassi traktorov. Pod obshchei red. D.K.Karel'akikh. 1950. 144 p. (MIRA 11:12) (Tractors)





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ZARFTSKIT, D. P. [Effective boundary conditions for "gray" bodies] Effektivnye granichnye unlovita dita "serykh" tel; doklady, predstavlennye SSSR na Meshdunarodnulu konferentsitu po mirnozu ispol'zovanitu atomnoi energii. Hoskva, 1955. 17 p. [Microfilm] (Nuclear physics) (MIRA 9:3)

USSR/Nuclear Physics - Thermal neutron capture

FD-3260

Card 1/1

Pub. 146 - 19/44

Author

: Shut'ko, A. V.; Zaretskiy, D. F.

Title

Capture of thermal neutrons by lead isotopes

Periodical

: Zhur. eksp. i teor. fiz., 29, No 6(12), Dec 1955, 867-868

Abstract

: The authors consider the isotopes Pb-207 and Pb-208 and their excitation levels, spins, parities, energies, etc. They compare the theoretical evaluations of cross-section of thermal neutron capture by lead isotopes with data of experiments. They conclude that capture in Pb-206 is "less single-particle" than in the case of Pb-207, and that the anomalous character of capture radiation in lead isotopes is explained by proceeding from the single-particle picture of capture. The authors thank Professor A. S. Davydov and V. F. Turchin for comments. Seven references, including one USSR: L. K.

Peker and L. A. Sliv, Izv. AN SSSR, ser. fiz., 17, 1953.

Institution

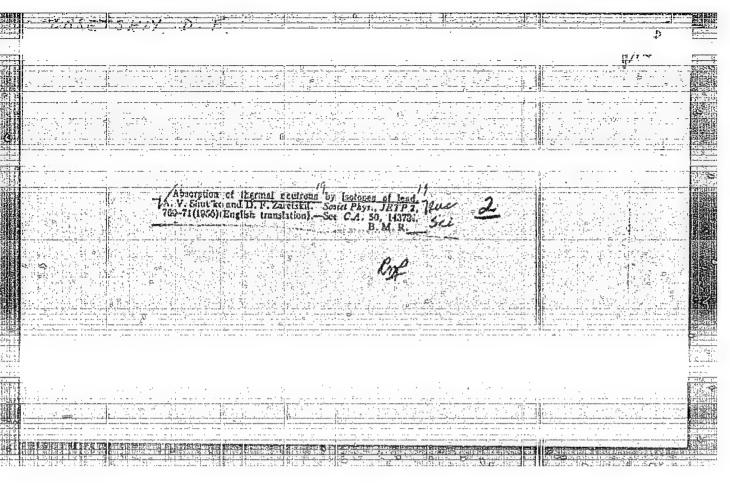
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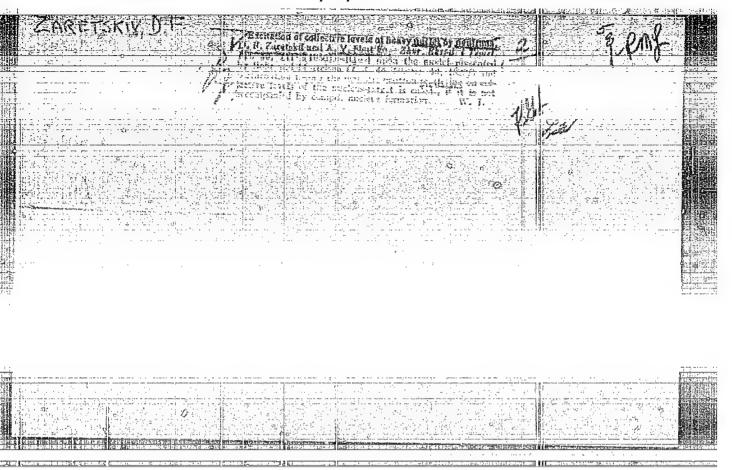
August 31, 1955

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"On the Thermal Capture by Fb Isotopes" a paper presented at the Enternational Conference on Nuclear Reactions, Amsterdam, 2-7 July 1956.

D551274





AUTHOR ZARETSKIY, D.F., SHUTKO, A.V. PA - 2690 On the Quasi-Magnetical Interaction of the Nucleon Spin With the TITLE Rotation of the Nucleus. (O kvazimagnitom vzaimodeystvii spina nukloma s vrashcheniyem yadra-Russian) Zhurmal Eksperim. i Teoret. Fiziki, 1957, Vol 32, Nr 2, pp 370-371, PERIODICAL (U.S.S.R.) Received 5/1957 Reviewed 6/1957 It is possible to find a new interpretation for the rotational levels of ABSTRACT' nuclei with spin 1/2 if we start out from the following premises: (1) In These nuclei there exist E-states. Then in first approximation the levels with the total angular momentum I=K±1/2 are degenerated. (Here K stands for the rotational quantum number). (2) This degeneration is eliminated if we introduce into the Hamiltonian by Bohr and Mottelson an interaction of the form of $H_{Rs} = (\lambda/mc^2)$ (70 $\forall_{koll})$ Here Astands for a mondimensional phenomenological constant with the same significance and magnitude as in the normal (usual) spin-orbit coupling of the nucleus. Furthermore the following denotations are used: I for the vector of the nucleon spin. $U(\vec{r})$ for the selfconsisting potential of the nucleus; me for the rest energy of the nucleon, and $v_{\rm koll}$ for the velocity with which the nucleon participates in the collective motion. First of all the sigmificance of vkoll is clarified, and then the significance and the ori-Card 1/2 gim of the above-mentioned interaction HRg. For the wave function we set

AUTHORS:

Bat', G.A., Zarotskiy, D.F.

304/19-4-6-2/30

TITLE:

The Effective Boundary Value Conditions in the Diffusion Theory of Neutrons (Surrey) (Effektivnyye granichnyye usloviya v teorii diffuzii neytronov (Obzor))

PERIODICAL:

Atomoaya energiya, 1958, Vol 4, Nr 6, pp 510-519 (USSR)

ABSTRACT:

In the course of a survey the methods are described by means of which the effective boundary value conditions can be determined which bring about agreement within the asymptotic range between the solution of the neutron diffusion equation and the solution of the corresponding kinetic equation.

The boundary values are described for monoenergetic neutrons; some of them are computed, and, especially, the occurring coefficients are registered in tables for the following cases:

1.) For an infinitely thin "black" rod in a medium, in which sources are uniformly distributed.

2.) For a plane boundary between medium and vacuum by means of the exact analytical method developed by Wiener and Hopf.

3.) For a "black" round cylinder. The following methods of solving are mentioned for this purpose:

Cerd 1/2

The Effective Boundary Value Conditions in the SOV/19-4-6-2/30 Diffusion Theory of Neutrons a) The balance method developed by Brudno (1951)
b) The variation method developed by Zaretskiy (Ref 6) o) The method of spherical harmonics by Galanin d) Solution of Peierl's (Payerl's) equations for great and small r, according to the method developed by Davison... 4.) For a "black" rod of any cross section by the method of approximation developed by Hurwitz (Gurvits) and Roe (Ref 8). 5.) For a "gray" rod (only possible by approximation). If it is invended to take the neutron spectrum into account in connection with the boundary value conditions, this is possible only for some simple border cases. (A body is described as being either black or gray according to whether the neutrons impinging upon its surface are fully or only partly absorbed). There are 2 figures, 5 tables and 10 references, 4 of which are Soviet. SUBMITTED: October 4, 1957 1. Neutrons--Diffrusion 2. Diffusion--Theory 3. Mathematics ---Applications Card 2/2

ZARETSKIY, D. F.

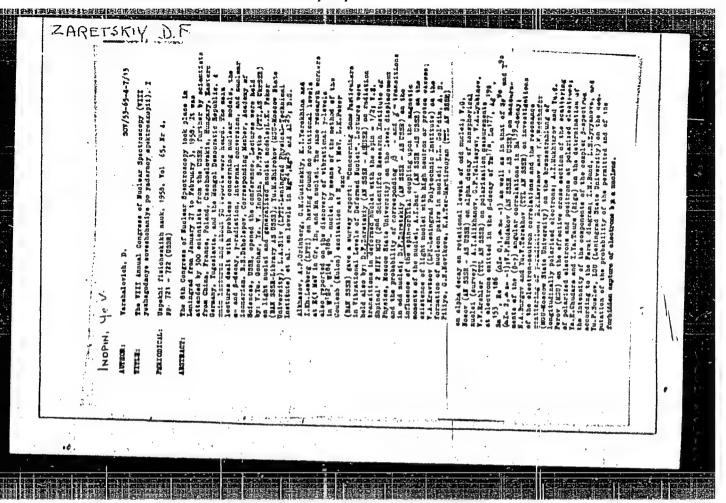
"On -Mesonic Fission."

paper to be presented at 2nd Un Intl. Conf. on the peaceful uses of Atomic Energy, Geneva, 1 - 13 Sep 58.

DROZDOV, S. I., ZARETSKIY, D. F., KUBRIN, L. P. and SEDELNIKOV, T. Kh.

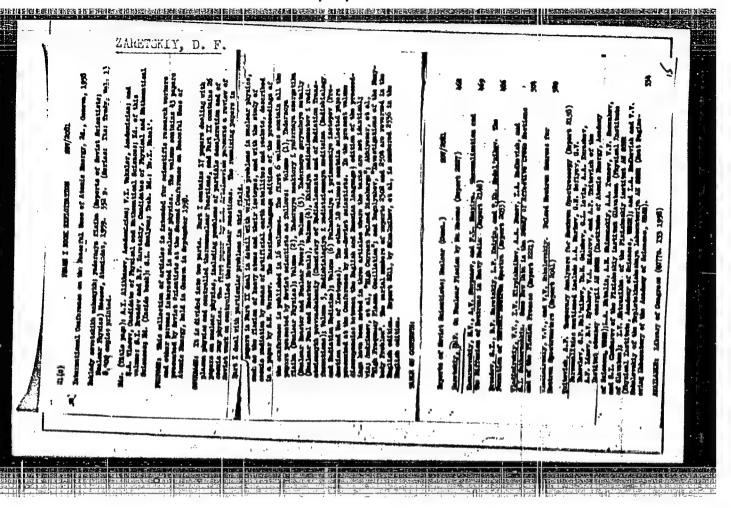
"On the Formation of a Thermal Neutron Spectrum."

Paper to be presented at 2nd UN Ints. Conf. on the peaceful uses of At aic Energy, Geneva, 1 - 13 Sep 58.



"APPROVED FOR RELEASE: 09/19/2001

CIA-RDP86-00513R001963820010-1



21(1),24(5) AULHOR: Zarotskiy. D. F. SOV/56-36-3-54/71 TITLU: On the Theory of Nuclear Magnetic Moments (K teorii magnitnykh momentov yader) . CHODICAL: Zhurnal okaperimental noy i teoreticheskoy fiziki. 1959, Vol 36, Nr 3, pp 869-873 (USSR) ADGY Adv: The important part played by spin-orbit coupling in the theory of nuclear shells is, in general, taken into account by (1) : $\chi'_{1a} = -(\lambda / \hbar/m^2c^2) \vec{s} [\nabla \vec{v}\vec{p}]$ λ = phenomenological spin-orbit coupling constant, λ = nucleon spin operator, λ = selfconsistent nuclear potential, λ = nomentum operator of the nucleon); if the nucleus is in an electromagnetic field, (1) with paca/c goes over into $\mathcal{X}' = (\lambda e \frac{\hbar}{m^2}c^3) \vec{s} \nabla u \vec{a}$ where a is the vector potential. Mayer and Jensen (Ref 1) were the first to point out the necessity of taking this spin-orbit effect into account without, however, using (2). Card 1/3

On the Theory of Ruclear Magnetic Moments

sov/56-56-5-54/71

D. P. Grochukhin (Ref 2) calculated the spin-orbit coupling for magnetic transitions of the kind $d_{3/2} \rightarrow s_{1/2}$. In the present paper the effect is investigated by which (2) is introduced into the magnetic moment of spherical and nonspherical nuclei with odd proton number. Although (2) is not the only cause of the deviation of the observed magnetic moments from the Schmidt (Shmidt) line, an investigation (estimate) of (2) is nevertheless of interest, because it facilitates separation of the other effects. In some cases (2) makes the main contribution to the deviation of magnetic moments from the Schmidt line as well as to the probability of forbidden M1-transitions for nuclei with an old proton number. In the case of nuclei with odd neutron numbers, consideration of (2) is of good purpose when there exists an effective charge due to interaction between noutron and the core of the nucleus (cf. Hef 3). In the present paper the additional factor occurring for strong spin-orbit coupling by the interaction between the nucleons in the nuclous and the electromagnetic field is, essentially, calculated for spherical and nonspherical nuclei. For the former the magnetic moment is $\mu = \langle J, J | \hat{\mu}_z | J, J \rangle$

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On the Theory of Ruelear Lagnotic Moments

sov/56-36-3-34/71

and, in consideration of (2):

$$\mu = \frac{1}{4} \frac{\frac{1}{2} \cdot \frac{1}{2}}{\frac{1}{2} \cdot \frac{1}{1} \cdot 1} \frac{e^{\frac{\lambda}{2}}}{2mc} \frac{\lambda}{mc^2} < J \left| r \frac{\partial U}{\partial r} \right| J >$$

For Ei 200 an estimate results in $\mu \approx +0.6$ e 200 e. For nonapherical nuclei it holds that $\mu = \langle \vec{\mu} \vec{J} \rangle/(J+1)$ and with (2) for spin 1/2 also a formula, (21), is derived. For Tm 109 this results in $\mu = -0.35$ nuclear magnetons, as against an experimental value of $\mu = -0.20 \pm 0.5$. The author finally thanks P. E. Nemirovskiy for discussions. There are 13 references, 7 of which are Soviet.

SUBMITTED:

September 16, 1958

Card 3/3

"APPROVED FOR RELEASE: 09/19/2001 C

CIA-RDP86-00513R001963820010-1

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24(5) AUTHOR:	Zaretskiy, D. F.	SOV/56-36-4-26/70		
TITLE:	On the Radiative Transitions in the System of Rotational Levels of Nuclei With Spin 1/2 (O radiatsionnykh perekhodakh v sisteme vrashchatel'nykh urovney yader so spinom 1/2)			
PERIODICAL:	Zhurnal eksperimental'noy i t Vol 36, Nr 4, pp 1129-1132 (U	eoreticheskoy fiziki, 1959, SSR)		
ABSTRACT:	For the purpose of calculating magnetic dipole- and electric between rotational levels of coupling system suggested by used by the author (Ref 1) is coordinates (transition to the nuclear axes) is carried and by using the relation X.	deformed spin 1/2 nuclei, a Hund and already previously sused. A transformation of ne system of coordinates of out in the wave function,		
Card 1/3		er .		

On the Radiative Transitions in the System of Rotational Levels of Nuclei With Spin 1/2

sov/56-36-4-26/70

a wave function of the form

$$\Psi_{xs}^{JH} = \sqrt{(2+1)/8\pi^{2}} y_{0} \sum_{m_{8}'} (x /_{2} 0 m_{8}' | x /_{2} J m_{8}') D_{M,m_{8}'}^{J} \chi_{\lambda,m_{8}'}$$

is obtained, which causes no difficulties with respect to calculating the probabilities of electric quadrupole transitions. The result agrees with that obtained by Bohr and Mottelson (Ref 2). (Yo denotes the wave function

describing the motion of the nucleon with respect to the nuclear axes, X- the spin function of an odd nucleon in the laboratory system, the rotation quantum number X

assumes the values 0,2,4.... for the states \sum_{g}^{+} and 1,3,5.... for \sum_{u}^{+} . In the following the authors

calculate the probabilities of magnetic dipole transitions and show that transitions between components of different doublets with $\Delta J=1$ occur mainly as magnetic dipole transitions. The admixture of electric quadrupole transitions amounts to not more than 10 - 20 %. In

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On the Radiative Transitions in the System of Rotational Levels of Nuclei With Spin 1/2

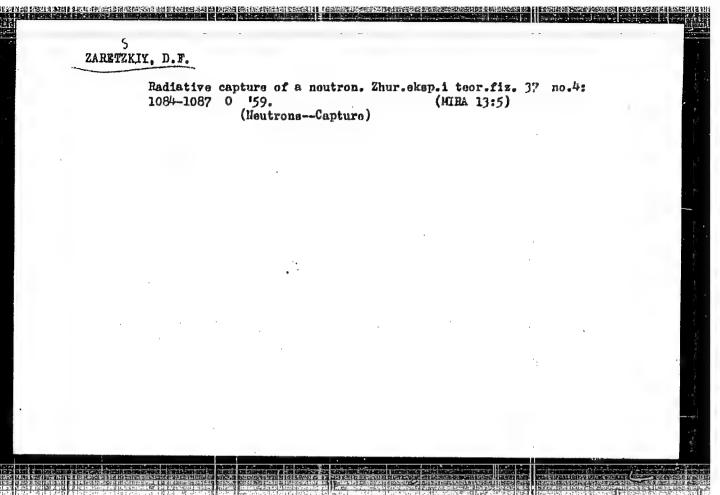
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conclusion, the nucleus Tm 169 (and Tm 171) is shortly investigated as an example, and a comparison is drawn with the results obtained by references 7 and 8 with respect to the nucleus Yb 171, the daughter nucleus originating from the \(\int_{\text{-decay}}^{\text{-decay}} \) of Tm 171 (Fig). The author thanks S. A. Baranov for discussing experimental data, and D. P. Grechukhin for his comments. There are 1 figure and 9 references, 3 of which are Soviet.

SUBMITTED:

September 13, 1958

Card 3/3



16.8100,24.6720,24.6800,24.6810

AUTHORS:

Zaretskiy, D. F., and Novikov, V. M.

TITLE:

Letter to the Editor. Depolarization of Muons in μ -Mesoatoms with Deformed Nuclei

PERIODICAL:

Zhurnal eksperimental noy i teoreticheskoy fiziki,

1959, Vol 37, Nr 6, pp 1824-1825 (USSR)

ABSTRACT:

An analysis was made of the additional depolarization caused by the interaction of muon with nuclear deformation in μ -mesoatoms. The effect was considered for even-even nuclei. The Hamiltonian of the muon-nucleus

system was taken in the following form:

 $H = H_0 + H_R + H_q$

Here, Ho is the Hamiltonian of the muon in the monopole field; Hp is operator of rotational energy of the

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Letter to the Editor. Depolarization of Muons in μ -Mescatoms with Deformed Nuclei

77009 SOV/56-37-6-49/55

nucleus; H_q is operator of quadrupole interaction of muon with nucleus. The diagonalization of the Hamiltonian shows that the quadrupole interaction considerably changes the eigen functions of the system with corresponding muon in the 2p-states (cf., L. Wilets, Kong. Dansk. Vidensk. Selsk. Mat.-fys. Medd., 29, 3, 1954). The changes of the eigen functions of other states can be neglected. In such an approximation the polarization of muon in 1s-states becomes:

 $P = A_q W_{V_s} \langle \sigma_{2\rho_{V_s}} \rangle_0 + B_q W_{V_s} \langle \sigma_{2\rho_{V_s}} \rangle_0$

Here $W_{\frac{1}{2}}$ and $W_{3/2}$ are probabilities of the passing of muon through states $2p_{\frac{1}{2}}$ and $2p_{3/2}$; $\langle \sigma_{2p_{\frac{1}{2}}} \rangle_o$ and

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Letter to the Editor. Depolarization of Muons in μ -Mesoatoms with Deformed Nuclei

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 $\sigma_{2p_{3/2}}$ o are polarization of muon in the above states in the absence of quadrupole interaction; Aq and Bq are factors compensative of the additional depolarization. The factor Bq was determined from muon transition from the upper levels into the ls-state through the states that can be described by the eigen function of the above given Hamiltonian:

$$B_q = 0.8 \sum_{k=1}^{2} C_h^1 \left[\left(\frac{E_{f,d} - 3E_h}{3E_q} \right)^2 + 2 \right] - 0.6,$$

$$C_{k}^{2} = \frac{E_{q}^{2} (E_{R} + E_{q} - \frac{2}{3} E_{f, \bullet} - E_{k})^{1}}{E_{q}^{2} (E_{R} + E_{q} - \frac{2}{3} E_{f, \bullet} - E_{k})^{2} + [(\frac{1}{3} E_{f, \bullet} - E_{k})^{2} + [(\frac{1}{3} E_{f, \bullet} - E_{k})^{2} + (\frac{1}{3} E_{f, \bullet} - E_{k})^{2}]}$$

$$\Rightarrow \frac{E_{q}^{2} (E_{R} + E_{q} - \frac{2}{3} E_{f, \bullet} - E_{k})^{2}}{-E_{k} (E_{R} - \frac{2}{3} E_{f, \bullet} - E_{k}) - E_{q}^{2}]^{2}}.$$

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Letter to the Editor. Depolarization of Muons in μ -Mesoatoms with Deformed

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Here, E_k is eigen values of the Hamiltonian which correspond to muon in 2p-states and to the whole momentum of the system 3/2; $E_{f.s.}$ is splitting of fine structure of the 2p-level; E_R is energy of the first rotational level of the nucleus; E_q = \langle 2p $|H_q|$ 2p \rangle . The following values of B_q and D_q were calculated for Gd 15d , W^{184} , Th^{232} , and U^{238} , respectively: 0.63, 0.64, 0.38, 0.38 and 0.61, 0.64, 0.6, 0.59. Thus, the interaction of muon with nuclear quadrupole deformation can lead to a considerable additional depolarization of the muon. There are 4 references; 2 Soviet, 1 Dutch, 1 U.S. The U.S. reference is: G. W. Ford, C. J. Mullin, Phys. Rev., 108, 477, 1957.

SUBMITTED:

August 21, 1959

Card 4/4

GRIN', Yu.T.; DROZOV, S.I.; ZARETSKIY, D.F.

Green's function for odd nuclei. Zhur. eksp. i teor. fiz. 38
no.1;222-228 Jan '60. (MIRA 14:9)

(Potential, Theory of) (Nuclei, Atomic)

8/048/61/025/009/005/007 B104/B102

Grin', Yu. T., and Zaretskiy, D. F.

TITLE:

Collective excitations of non-spherical nuclei

PERIODICAL: Akademiya nauk SSSR. Izvestiya. Seriya fizicheskaya, v. 25, no. 9, 1961, 1169 - 1175

TEXT: This paper was read at the 9th Annual Conference Nuclear Spectroscopy. The authors generalize the theory of collective excitations of non-spherical nuclei. An equation is set up for the frequencies of collective nuclear oscillations, and a relation between the frequencies of the excited $\beta(k=0)$ and f'(k=2) vibrational levels is derived in quasiclassical approximation. Using the Green's two-particle function

 $K = \langle \Phi_0 | T(a_1 a_2^+ a_3 a_4^+) | \Phi_0 \rangle$, the equation

 $1 = \varkappa \sum_{12} \frac{(E_1 E_1 - \varepsilon_1 \varepsilon_2 + \Delta^2)}{2E_1 E_2 \left[\omega^2 - (E_1 + E_1)^2 \right] \left| (q_{2k})_{12} \right|^2}$ (9)

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Collective excitations of non- ...

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is obtained for the frequencies of the bound states of two quasi-particles. Within the framework of the Nilsson model Eq. (9) can only be solved numerically for a real model. Using the model of an axisymmetrically deformed oscillator potential, the authors attempt to find a solution to this equation. Eq. (9) is represented in the form

$$1 = \times \sum_{\lambda \lambda'} \frac{2\Delta^2 |q_{\lambda \lambda'}|^3}{E_{\lambda} (4E_{\lambda}^2 - \omega^2)} + \times \sum_{\lambda \lambda'} \frac{(E_{\lambda}E_{\lambda'} - e_{\lambda}e_{\lambda'} - \Delta^2) (E_{\lambda} + E_{\lambda'})}{2E_{\lambda}E_{\lambda'} [(E_{\lambda} + E_{\lambda'})^3 - \omega^2]} |q_{\lambda \lambda'}|^2, \quad (15)$$

The first term is the sum over all transitions without energy change, and the second term is the sum over all other transitions. Since () is negligible in the latter sum, Eq. (15) can be reduced to

$$1 = \varkappa' \sum_{\lambda\lambda'} \frac{2\Delta^a}{B_\lambda \left(4E_\lambda^2 - \omega^a\right)} |q_{\lambda\lambda'}|^a,$$

where

$$\kappa' = \frac{\kappa}{1 - \kappa \sum_{\lambda \lambda'} \frac{E_{\lambda} E_{\lambda'} - \varepsilon_{\lambda} \varepsilon_{\lambda'} + \Delta^{2}}{2E_{\lambda} E_{\lambda'} (E_{\lambda} + E_{\lambda'})} |q_{\lambda \lambda'}|^{2}}$$
(16)

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Collective excitations of non- ...

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In quasi-classical approximation, Eq. (16) leads to

$$\frac{1-\lambda^{1}-\frac{4\Delta^{2}}{\omega^{2}-1}}{\omega^{2}\sqrt{\frac{4\Delta^{2}}{\omega^{2}-1}}} = \frac{1}{\sqrt{\frac{4\Delta^{2}}{\omega^{2}-1}}} = \frac{1}{\lambda \lambda^{1}} \left| \frac{2}{\delta(\epsilon_{\lambda})} \right| \lambda \sim \epsilon_{0}/AR_{0}^{4}.$$

The following explicit solutions are obtained from this equation: $\omega = 2\Delta \sqrt{\frac{3}{2}} \cdot \sqrt{1-z'} \frac{\sum_{\lambda \lambda'} |q_{\lambda \lambda'}|^2 \delta(\epsilon_{\lambda})}{\sqrt{1-z'}}$ for $\omega \ll 2\Delta$ and

 $\omega = 2\Delta \left\{ 1 - \frac{\pi}{8} z^{1/2} \left[\sum_{\lambda \lambda'} \left| q_{\lambda \lambda'} \right|^2 \delta(\epsilon_{\lambda}) \right]^2 \right\} \qquad \text{for } 2\Delta - \omega \ll 2\Delta.$

An analysis of these solutions indicates that collective excitations with energies much less than 2Δ may exist in deformed nuclei. β - and β -

Card 3/5

27480 \$/048/61/025/009/005/007 B104/B102

Collective excitations of non-

$$\frac{\omega_{\gamma}^{2}}{\omega_{\gamma}^{2}} = \frac{\sum_{\lambda\lambda'}^{1} \frac{1 - g\left(\frac{\varepsilon_{\lambda} - \varepsilon_{\lambda'}}{2\Delta}\right)}{\left(\frac{\varepsilon_{\lambda} - \varepsilon_{\lambda'}}{2\Delta}\right)^{3}} (q_{22})_{\lambda\lambda'}^{2} \delta(\varepsilon_{\lambda})}{\sum_{\lambda\lambda'}^{1} \frac{1 - g\left(\frac{\varepsilon_{\lambda} - \varepsilon_{\lambda'}}{2\Delta}\right)}{\left(\frac{\varepsilon_{\lambda} - \varepsilon_{\lambda'}}{2\Delta}\right)^{3}} (q_{30})_{\lambda\lambda'}^{3} \delta(\varepsilon_{\lambda})}.$$
(23)

From this relation, it is concluded that the frequencies of β - and γ -vibrations agree within the framework of the model of an axisymmetric oscillator potential. The agreement of the frequencies is closely related to the degeneracy of the levels in the oscillator potential. Finally, the Nilsson model without spin-orbital coupling is briefly discussed. The following values are obtained for the ratio $\alpha_{\beta}/\alpha_{\gamma}$ as a function of

 ω_{0}/ω_{1} 1.8 2.25 3.0 4.5 ω_{0}/ω_{1} 0.90 0.86 0.80 0.72

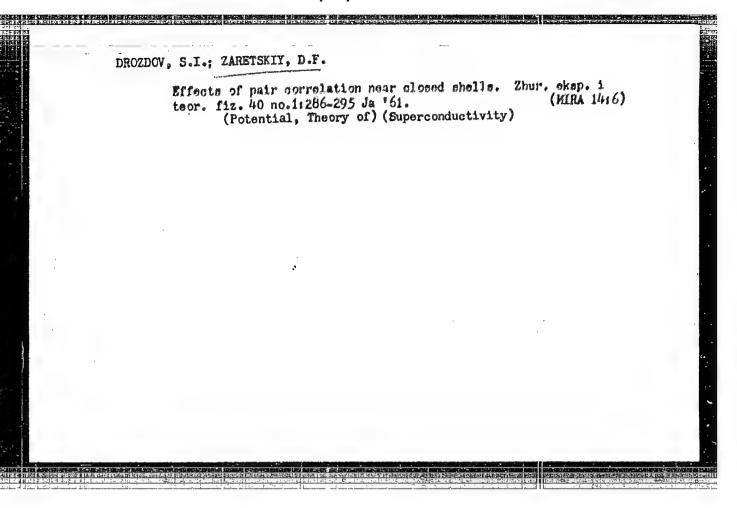
S. T. Belyayev (Zh. eksperim. i teor. fiz. 39, 1387 (1960)) is mentioned. Card 4/5

27\80 8/048/61/025/009/005/007 B104/B102

Collective excitations of non- ...

There are 11 references: 7 Soviet and 4 non-Soviet. The references to English-language publications read as follows: Marumori T., Progr. Theor. Phys., 24, 351 (1960); Baranger M., Phys. Rev., 120, 957 (1960); Perlman I., Proceedings of the International Conference on Nuclear Structure, p. 547, Kingston, Canada, 1960.

Card 5/5



S/056/61/040/003/031/031 B112/B214

24.6900 AUTHORS:

Zaretskiy, D. F., Novikov, V. M.

TITLE:

Nuclear fission by μ -mesons

PERIODICAL:

Zhurnal eksperimentalinoy i teoreticheskoy fiziki, v. 40, no. 3, 1961, 982-983

TEXT: The effect of a meson on nuclear fission can be calculated for the case of rigidly oriented nuclear axes. The energy \mathbf{E}_μ of a bound meson

depends on the deformation parameters of the nucleus. A solution of Schrödinger's equation for a meson in the Coulomb field of the deformed nucleus is required for the determination of \mathbf{E}_{μ} . It is assumed in the

present paper that the nucleus has the form of an ellipsoid of rotation up to its saddle point. The Coulomb potential of a homogeneously charged ellipsoid of rotation with the semiaxes a and b has the form:

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S/056/61/040/003/031/031 B112/B214

Nuclear fission by ...

$$\varphi(a, \beta) = \frac{Ze}{c} \left\{ [1 - P_3(\cosh \alpha) P_2(\cos \beta)] \ln \coth \frac{\alpha_0}{2} + \frac{3}{2} \frac{\cosh^2 \alpha}{\cosh \alpha_0} P_2(\cos \beta) + \frac{3}{4} \left(1 - \frac{\sinh^2 \alpha}{\sinh^2 \alpha_0} \right) \frac{\sin^3 \beta}{\cosh \alpha_0} \right\} \quad \text{for } \cosh \alpha \leqslant \cosh \alpha_0 = \frac{a}{c},$$

$$\varphi(a, \beta) = \frac{Ze}{c} \left\{ [1 - P_2(\cosh \alpha) P_3(\cos \beta)] \ln \coth \frac{\alpha}{2} + \frac{3}{2} \cosh \alpha P_3(\cos \beta) \right\}$$
(1)

for $\operatorname{ch} \alpha \geqslant \frac{a}{c}$,

Here, Ze is the nuclear charge, $c^2 = a^2 - b^2$, P is the Legendre polynomial, and α,β are the degenerate elliptical coordinates. The Schrödinger equation for a bound meson in a nuclear field with such a potential was solved numerically with the help of an electronic computer. The value of E_{thresh} for different values of a/b is given in the following table (U^{238} nucleus):

following table (U^{238} nucleus): a/b = 1.2 1.4 1.6 1.8 2 2.2 2.5 E. (Mey) = 11.89 11.78 11.66 11.53 11.36 11.21 11.0

The increase AE of the fission barrier for some nuclei is given in another table:

Card 2/3

Nuclear fission by ...

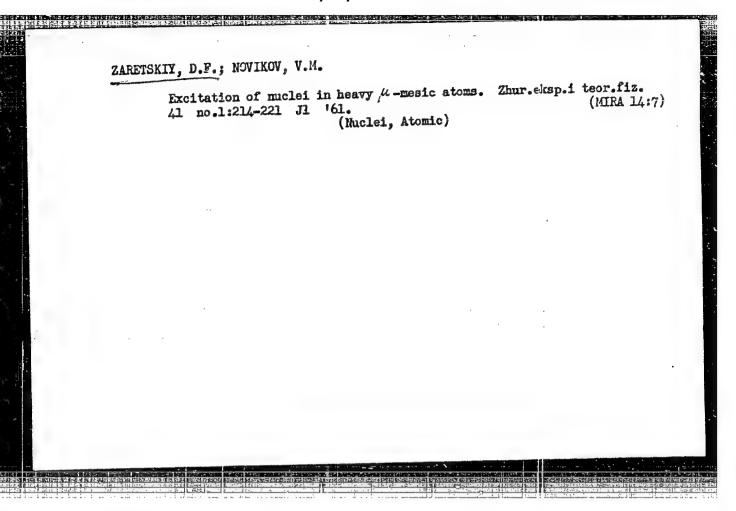
S/056/61/040/003/031/031 B112/B214

Nucleus	a/b statistical	a/b saddle point	Ethreshold, Mev	AE, Mev
ປ ² 3ອີ	1.30	2.24	5.8	0.6
ປ ² 3ອີ	1.25	2.2	5.75	0.6
Pu ² 39	1.30	2.17	5.48	0.5

D. P. Grechukhin is thanked for his advice and V. K. Saul'yev for setting up the program and calculation with the electronic computer. There are 2 tables and 7 references: 2 Soviet-bloc and 3 non-Soviet-bloc.

SUBMITTED: January 11, 1961

Card 3/3



ZARETSKIY, D.F.; URIN, M.G.

Nature of collective levels in nonspherical nuclei. Zhur.eksp.
i teor.fiz. 41 no.3:898-906 S '61. (MIRA 14:10)

1. Moskovskiy inzhenerno-fizioheskiy izstitut.
(Nuclei, Atomic)

ZARETSKIY, D.F.; JIRIN, M.G.

Microscopic description of collective levels of nonspherical nuclei.
Zhur. eksp. i teor. fiz. 43 no.3:1021-1030 '62. (MRa 15:10)

1. Moskovskiy inzhonerno-fizicheskiy institut.
(Quantum theory) (Nuclie, Atomic)

"APPROVED FOR RELEASE: 09/19/2001 CIA-RI

s/056/62/042/001/045/048 B154/B112

AUTHORS:

Zaretskiy, D. F., Urin, M. G.

TITLE:

Low levels with negative parity in nonspherical nuclei

PERIODICAL:

Zhurnal eksperimental'noy i teoreticheskoy fiziki, v. 42,

no. 1, 1962, 304-305

TEXT: In deformed even-even nuclei, low levels with 0.3-1.3 Mev were recently detected (Ref. 1). These 1- levels correspond to internal excitations with K = 0 and negative parity. They cannot be explained by the hydrodynamic model. Therefore, it was proposed that such levels should be interpreted in terms of the superfluid model. Using the method of approximate second quantization in the same way as in Ref. 1 (Zaretskiy, D. F., Urin, M. G., ZhETF, 41, 398, 1961), the excitation energy (N = 1) is found to be

1s found to be $1 = x \sum_{\lambda \lambda^{\dagger}} n_{,p} |(q_{30})_{\lambda \lambda^{\dagger}}|^{2} \frac{E_{\lambda}E_{\lambda^{\dagger}} - E_{\lambda}E_{\lambda^{\dagger}} + \Delta^{2}}{2E_{\lambda}E_{\lambda^{\dagger}}} \frac{E_{\lambda} + E_{\lambda^{\dagger}}}{(E_{\lambda} + E_{\lambda^{\dagger}})^{2} - \omega^{2}}$ (1)

Card 1/5

s/056/62/042/001/046/048 B154/B112

Low levels with negative ...

where $q_{30} = r^3 \cdot \gamma_{30}(\theta)$ is the single operator for the octupole moment; $E_{\lambda} = \sqrt{\mathcal{E}_{\lambda}^2 + \Delta^2}$. The energy \mathcal{E}_{λ} is substracted from the Fermi surface; $\Delta = \Delta_n(\Delta_p)$ is a constant characterizing the energy of proton (neutron) pair correlation; $\mathcal{E}_0^n = \mathcal{E}_0^n(\mathcal{E}_0^p)$ is the energy of the Fermi boundary for neutrons (protons); $\kappa \sim \mathcal{E}_0/AR^6$ is constant depending on the octupole-octupole interaction. The value of κ is proposed to be the same for nn, pp, and np interactions. (1) permits two solutions which, at $\kappa \to 0$, correspond to the dissociation of a neutron or a proton pair. For a given group of nuclei, κ can be regarded as constant. If κ is therefore determined for one group of nuclei from the position of the 1 level, it will be possible to estimate ω from Eq. (1) for the other nuclei of this group. The probability of electric dipole transition from the 1 level to the ground level is expressed by

Card 2/5

Low levels with negative ...

\$/056/62/042/001/046/048 B154/B112

$$B(E1; I^{-} \rightarrow 0^{+}) =$$

$$= \frac{1}{6\omega} \left\{ \sum_{\lambda\lambda'} e_{n, p} \left(q_{10} \right)_{\lambda\lambda'} \left(q_{30} \right)_{\lambda\lambda'} \frac{E_{\lambda}E_{\lambda'} - \epsilon_{\lambda}e_{\lambda'} + \Delta^{2}}{2E_{\lambda}E_{\lambda'}} \frac{E_{\lambda} + E_{\lambda'}}{(E_{\lambda} + E_{\lambda'})^{3} - \omega^{3}} \right\}^{2} \times \left\{ \sum_{\lambda,\lambda'} e_{n, p} \left| \left(q_{30} \right)_{\lambda\lambda'} \right|^{2} \frac{E_{\lambda}E_{\lambda'} - \epsilon_{\lambda}\epsilon_{\lambda'} + \Delta^{2}}{2E_{\lambda}E_{\lambda'}} \frac{E_{\lambda} + E_{\lambda'}}{\left[(E_{\lambda} + E_{\lambda'})^{3} - \omega^{3} \right]^{3}} \right\}^{-1} \right\}$$

$$(2),$$

where $e_n=-Ze/\Lambda$, $e_p=Ne/\Lambda$, e is the proton charge, and Z(N) is the number of protons (neutrons). In quasiclassical approximation, the authors estimate B(E1) for low 1 levels $(\omega \ll 2\Delta_{n,p})$ according to (Ref. 4):

$$\sum_{\lambda,\lambda'} |(q_{30})_{\lambda\lambda'}|^2 \frac{E_{\lambda}E_{\lambda'} - \varepsilon_{\lambda}\varepsilon_{\lambda'} + \Delta^2}{2E_{\lambda}E_{\lambda'}(E_{\lambda} + E_{\lambda'})^3} \approx \frac{1}{4\Delta^2} \sum_{\lambda\lambda'} |(q_{30})_{\lambda\lambda'}|^2 \varphi\left(\frac{\varepsilon_{\lambda} - \varepsilon_{\lambda'}}{2\Delta}\right) \delta(\varepsilon_{\lambda}) \sim \rho_0 R^4 \Delta^{-2},$$
(3),

where $\varphi(x) = x^{-2} - \ln(x + \sqrt{1 + x^2}) \cdot \left[x^3\sqrt{1 + x^2}\right]^{-1}$, and φ_0 is the level energy $\sqrt{\frac{1}{2}}$

Low levels with negative ...

S/056/62/042/001/046/048 B154/B112

density near the Fermi surface. For the approximation in (3), $\xi_{\lambda} - \xi_{\lambda}$, $\lesssim 2\Delta$ is necessary. From the Nilson scheme it follows that in the observed nuclear quadrupole deformations β_0 there are levels λ , λ' for which this condition is satisfied. Analogous evaluation of the numerator in Eq. (2) leads to

 $B(E1, 1^- \to 0^+) \sim \left(\frac{N-Z}{A}\right)^2 (eR)^2 \beta_0^2 \rho_0 \Delta \frac{2\Delta}{\alpha i}$ (4).

In the model considered, the probability for the excitation of the 3 level is

$$B(E3, 0^{7} \to 3^{7}) = \frac{1}{2\omega} \left\{ \sum_{\lambda\lambda'} |(q_{30})_{\lambda\lambda'}|^{2} \frac{E_{\lambda}E_{\lambda'} - e_{\lambda}e_{\lambda'} + \Delta^{2}}{2E_{\lambda}E_{\lambda'}} \frac{E_{\lambda} + E_{\lambda'}}{(E_{\lambda} + E_{\lambda'})^{3} - \omega^{4}} \right\}^{1} \times \left\{ \sum_{\lambda\lambda'} |(q_{30})_{\lambda\lambda'}|^{2} \frac{E_{\lambda}E_{\lambda'} - e_{\lambda}e_{\lambda'} + \Delta^{2}}{2E_{\lambda}E_{\lambda'}} \frac{E_{\lambda} + E_{\lambda'}}{((E_{\lambda} + E_{\lambda'})^{3} - \omega^{4})^{3}} \right\}^{-1}.$$
(5)

This relation is obtained without introduction of the effective nucleonic charge for the E3 transition. Estimating B(E3, $0^+ \rightarrow 3^-$) in analogy to (3) Card 4/5

S/056/62/042/002/032/055 B108/B104

AUTHORS: Zaretskiy, D. F., Novikov, V. M.

TITLE: Excitation of collective nuclear levels in heavy pomesic

atome

PERIODICAL: Zhurnal eksperimental'noy i teoreticheskoy fiziki, v. 42, no. 2, 1962, 511 - 519

TEXT: The excitation probability of low collective levels with spin 2th which depends on the magnitude of the nuclear quadrupole moment and on the quadrupole transition probability for the level in question is discussed for heavy mesic atoms. Not only a passing 2p muon but also a 3d muon may excite the rotational levels of heavy Th or U type nuclei. The interaction part of the Hamiltonian of the nucleus-muon system is ascribed to quadrupole interaction of the levels 2p_{1/2} and 2p_{3/2} only. Results of the calculations given in Table 1 indicate that the excitation probability W depends on the nature of the collective level. W is increased, 2th however, by some 20% owing to 3d_{3/2} and 3d_{5/2} muon states. The latter can Card 1/3

Excitation of collective nuclear ...

\$/056/62/042/002/032/055 \$108/B104

also lead to the excitation of the 4^+ level. The excitation probability $\frac{W}{4}$, however, is much more sensitive to the sign of the quadrupole moment

than is \mathbb{W}_{2}^+ . The hyperfine splitting of the excited nuclear levels owing

to the interaction between the magnetic moments of the muon and of the excited nucleus is calculated. The finite size of the nucleus has to be considered. For Th²³² and U²³⁸, $\Delta E = 840$ ev. The additional depolarization of muons owing to the transfer of polarity to a nucleus with excited collective levels is also discussed. However, this effect may be masked by the effect of the electron shell. The circular polarization of gammas from the 2p - is transition may give information about the additional depolarization of the muons. The effects of quadrupole interaction may also have an effect upon the shape of the muon transition lines. In general, the effect of quadrupole excitation of collective levels in heavy mesic atoms may be masked by radiationless excitation. The latter effect may be

estimated by studying the nuclear transitions $4^+ \rightarrow 2^+$ and $2^+ \rightarrow 0^+$ with respect to the occurring gamma quanta from the muon transition. There are 1 figure, 4 tables, and 13 references: 8 Soviet and 5 non-Boviet.

Card 2/3

Excitation of collective nuclear...

S/056/62/042/002/032/055 B108/B104

The four references to English-language publications Feedows 15R001963820010-1"

Jacobson. Phys. Rev. Lase 199/19/19001 M. GIAPROPES6-00515R001963820010-1"

42,APPROVED FOR RELEASE 199/19/2001 M. GIAPROPES6-Nucl. Phys., 14,
Nucl. Phys., 7, 569, 1958.

SUBMITTED: August 8, 1961

Table 1. Results.

Legend: (1) nucleus; (2) parameter of axial symmetry; (3) spin-orbital splitting; (4) energy of the 2+ level; (5) excitation probability.

14(7)	В (ЕЗ) -104 см ^а ОнЗ	2) Y, spad	B (13)	AB, (3)	Egs. May	w.C
Os180 Hg188	356 255 113	15,8 22,3 24,3	3,86 2,66 1,90	0,134 0,147 0,162	0,123 0,187 0,411	0,40° 0,30

Card 3/3

S/056/62/043/003/043/063 B108/B102

24.4400 AUTHORS:

Zaretskiy, D. F., Urin, M. G.

TITLE

|Microscopic description of collective levels of nonspherical |nuclei

PERIODICAL

Zhurnal eksperimental noy i teoreticheskoy fiziki, v. 43, no. 3(9), 1962, 1021 - 1030

TEXT: When pair correlations of nucleons are considered by the method of Hogolyubov's canonical transformation this leads to unphysical states. Here the structure of such collective states in deformed nuclei is studied by second quantization approximation (cf. ZhETF, 41, 898, 1961). In this case the Hamiltonian for the collective excitations is represented in a diagonalized form wherefrom the collective unphysical state can be eliminated. The unphysical states of the form $A_{11}^+|C_0\rangle$ are orthogonal to all physical (real) states. The error which results from simplifying the Hamiltonian $H = H_0 + H_0 + H_{int}$ to collective oscillations is of the order of $(Q_0\Delta)^{-1} \ll 1$ where Δ is a constant determining the pairing energy Card 1/3

Microscopic description of ...

S/056/62/043/003/043/063 B108/B102

and ϱ_0 is the energy density of the single-nucleon levels near the Fermi level. The matrix element of the electrical multipole moment for a transition between ground state and excited state with multipolarity λ and projection K of the moment has the form,

$$\langle \lambda K | Q_{\lambda K}^{el} | 0 \rangle^{2} = \left[e \sum_{1,2} | (q_{\lambda K})_{12} |^{2} I_{13} (\omega_{\lambda K}) \right]^{3} \left\{ \omega_{\lambda K} \sum_{1,2,p} | (q_{\lambda K})_{12} |^{2} J_{13} (\omega_{\lambda K}) \right\}^{-1}, \tag{24}$$

with $J_{12}(\omega^2) = \partial I_{12}/\partial \omega^2$; $I_{12}(\omega_{\lambda K}) = (E_1 E_2 - \ell_1 \ell_2 + \Delta^2) E_{12}/2 E_1 E_2 (E_{12}^2 - \omega_{\lambda K}^2)$; $E_{\nu} = \sqrt{\epsilon_{\nu}^2 + \Delta^2}$, $E_{12} = E_1 + E_2$, ϵ_{ν} is the single-nucleon state energy. In the case of β -excitations, Eq. (24) determines the reduced $2E\lambda$ -transition probability $B(E\lambda) \sim B_{\rm sp}(E\lambda) \rho_0 \Delta 2\Delta/\omega_{\lambda K}$ with an accuracy of $\Delta \beta_0$ where β_0 is the equilibrium deformation of the nucleus and $B_{\rm sp}$ is the reduced single-particle transition probability.

Card 2/3

ZARETRIKIY, D. F.; LUSHIHKOV, A. A.

"Applications of the Theory of the Fermi Liquid to Nuclear Photoabsorption."

report submitted for All-Union on Nuclear Spectroscopy, Toilisi, 14-22 Feb
64.

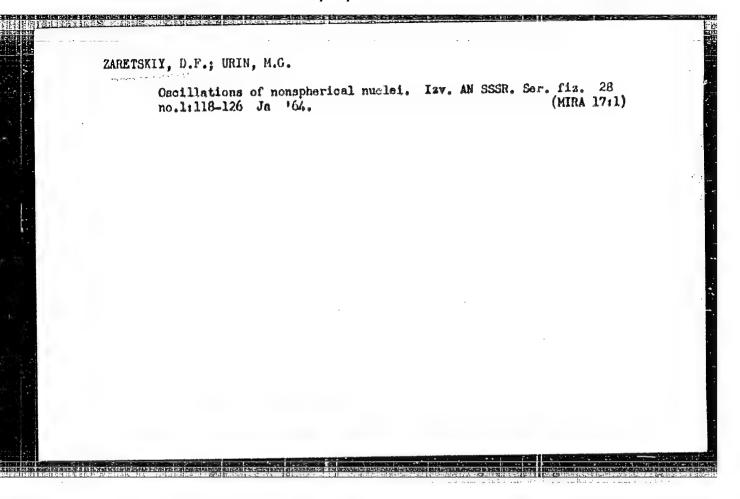
IAE (Inst Atomic Energy)

MIGDAL, A. B.; LUSHNIKOV, A. A.; ZARETSKIY, D. F.

"Theory of nuclear dipole photoabsorption."

report submitted for Intl Conf on Low & Medium Energies Nuclear Physics, Paris, 2-8 Jul 64.

Kurchatov Inst, Moscow.



ABSTRACT: The conditions are investigated, under which a gazze quantum radiated by one of the nuclei without recoil (Mossbauer effect) can be insorbed by another nucleus of the same type, so that the excitation (nuclear exciton) can propagate over an entire host crystal consisting of excited and unexcited nuclei of the same type. It is shown that when a nuclear exciton is produced, the main characteristics of spontaneous emission change markedly, because square anisotropy appears and the gazza quantum fluit is concentrated predominantly in this direction of the reciprocal-lattice vector. The emission probability in this direction is

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ACCESSION NR: AF5004413

nuclei. It is also shown that the formation of a nuclear exciton is connected with an increase in the width of the emitting level by a factor proportional to the cube root of the number of effective nuclei. It is thus preside to obtain sharply directional beams of monochromatic gamma quanta and the lifetimes of the nuclear increase can be erectly reduced when such isomers are placed in a crystal consisting of unexcited nuclei of the same sort. We thank F. .. Shapiro for wellable remarks and also V. K. Voytovelskiy and 8. M. Feynber: for discussions. Orig. ant. hes: 30 formalies.

ASSOCIATION: Institut atomicy energii (Insititute of Atomic Fuergy)

SCHREFFED: 28Julé4

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Card 2/2

"APPROVED FOR RELEASE: 09/19/2001

CIA-RDP86-00513R001963820010-1

L 2752-66 EWI(m)/I/EWA(m)-2
ACCESSION NR: AF5024345

AUTHOR: Zaretakiy, D. F.; Ivanter, I. G.

TITLE: The four-fermion interaction and baryon masses

SOURCE: Yadernaya fizika, v. 2. no. 2, 1965, 307-314

TOPIC TAGS: particle symmetry, unitary symmetry, group theory, strong nuclear interaction, baryon, fermion, particle physics

ABSTRACT: A simple dynamic model is proposed for interpreting the mass spectrum of baryons (octet and decuplet). The mathematical analysis is based on a Hamiltonian with strong four-fermion interaction, which corresponds to a scalar and to the eighth component of an 8-vector SU(3) representation. The four-fermion interaction is made up of strongly interacting baryon currents. This model is used as a basis for determining the relationship between splitting of the massis of the baryons in the octet and corresponding splitting in the decuplet. Orig. art. has: 32 formulas.

ASSOCIATION: none SUBMITTED: 30Dec64 NO REF SOV: 002

ENCL: 00 OTHER: 016

SUB CODE: HP, HA

SHAPIRO, I.S.; ZARETSKIY, D.F.; LUSHNIKOV, A.A.

[Nuclear physics; the mechanism underlying muclear reactions] IAdernaia fizika; mekhanizm iadernykh reaktsii. Moskva, AN SSSR, 1965. 88 p. (MIRA 18:10)

1. Akademiya nauk SSSR. Institut nauchnoy informatsii.

AUTHOR: Zaretskiy. D. F.; Lomonossy, Y. Y. ORG: none TITLE: Concerning certain features of the radiation of gaussa quanta from nuclei in a crystal lattice SOURCE: Yadernaya fizika, v. 3, no. 2, 1966, 263-267 SOURCE: Yadernaya fizika, v. 3, no. 2, 1966, 263-267 TOPIC TAGS: gaussa radiation, excited nucleus, nuclear isomer, nucleir energy level, crystal lattice structure, line width, spectral distribution, angular distribution ADSTRACT: This is a continuation of earlier work by the authors (ZhETF v. \$3, 568, crystal lattice structure, line width, spectral distribution collective nuclear excitation answers in an ideal single crystal, and that this exciton has a (nuclear exciton) can exist in an ideal single crystal, width of the nuclear level. The (nuclear exciton) can exist in an ideal single crystal width of the nuclear level of nuclear exciton decay width which differs noticeably from the natural width of radiative decay of nuclear exciten article is devoted to a study of the singularities of radiative decay citon lear levels in the case when the conditions for the excitence of a nuclear exciton are present, and to a study of this end, one of the methods of generation of nuclear excitention and the effects associated with it is investigated, and the quantum flux in this case. To this end, one of the methods of generation of nuclear radiation from a recoilless nucleus placed in a definite crystal littice point is decollective excitation and the effects associated with it is investigated, and the termined. This recoilless nucleus is assumed for simplicity to be a two-level isomer. Radiation from a recoilless nucleus is assumed for simplicity to be a two-level isomer. Cord 1/2	"APPROVED FOR RELEASE: 09/19/2001 CIA-RDP86-00	513R001963820010-1
	AUTHOR: Zaretskiy, D. F.; Lomonogoy, V. V. ORG: none TITLE: Concerning certain features of the radiation of gamma quanta crystal lattice SOURCE: Yadernaya fizika, V. 3, no. 2, 1966, 263-267 SOURCE: Yadernaya fizika, v. 3, no. 2, 1966, 263-267 TOPIC TAGS: gamma radiation, excited nucleus, nuclear isomer, nucleary stal lattice structure, line width, spectral distribution, angule crystal lattice structure, line width, spectral distribution, angule crystal lattice structure, line width, spectral distribution of earlier work by the authors (7 and 1965) where it is shown that under certain conditions collective in 1965) where it is shown that under certain conditions collective in (nuclear exciton) can exist in an ideal single crystal, and that the (nuclear exciton) can exist in an ideal single crystal, and that the case when the conditions for the existence of a present article is devoted to a study of the singularities of radial lear levels in the case when the conditions for the existence of the present article is devoted to a study of the spectral and angular distributions are present, and to a study of the spectral and angular distributions are present, and to a study of the spectral with it is investigated.	ear energy level, ar distribution there v. 43, 568, aclear excitation has a nuclear level. The ative decay of nuclear exciton ion of the gamma-neration of nuclear tigated, and the lattice point is de-
第二十二章 1915年1915年 2月 2月 2月 2日 2月 2日		

L 36377-66 nucleus, and line splitting is neglected. From the equations of motion the authors determine the amplitude of the state corresponding to the case when the isomer is excited but the other nuclei are not, the time dependence of the decay probability of the isomer, and the spectral and angular distribution of the gamma quanta emitted under definite relations between the wavelength of the radiation and the lattice constants. The results are extended to a three-level isomer and to radiation of more than one nucleus in the crystal. It is concluded that the existence of the nuclear exciton can be ascertained by observation of the spectral and angular distribution of the gamma-quantum flux when the condition $K = 2\pi b$ is satisfied (K - wave vector of the radiation, b - reciprocal lattice vector). The authors thank Y. K. Yoyto-vetskiy for useful discussions. Orig. art. has: 16 formulas. OTH REF: 001 ORIG REF: 002/ SUBM DATE: 08 Jan65/ SUB CODE: 20/ 2/2

"APPROVED FOR RELEASE: 09/19/2001 CIA-

CIA-RDP86-00513R001963820010-1

The effect of deformation on the corresion of metals. E.M. Zaratskii. J. Applied Chem. U.S. S.R. 24, 521-8(1951)(Engl. translation).

See C.A. 46, 4924a.

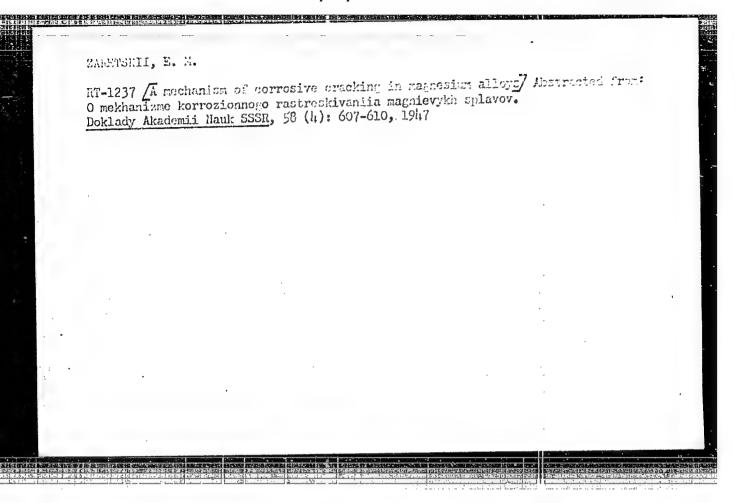
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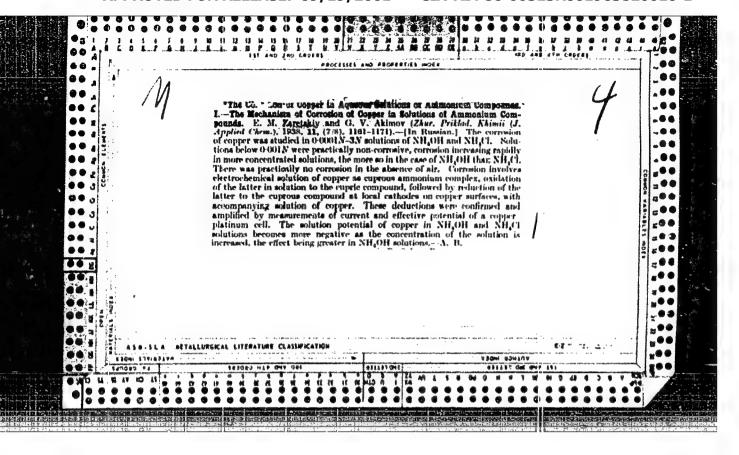
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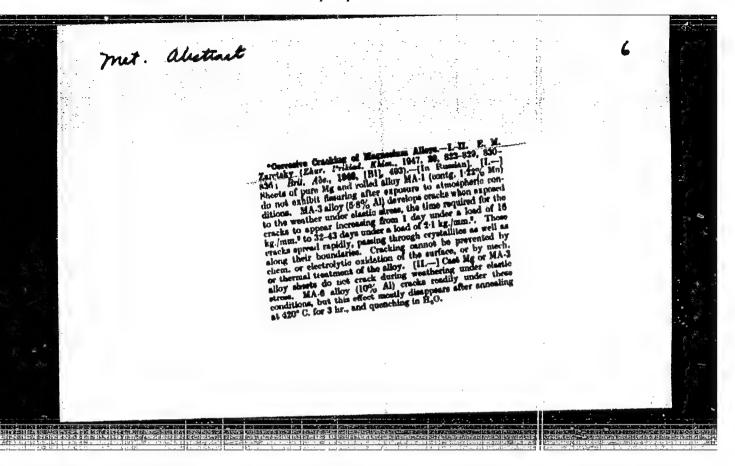
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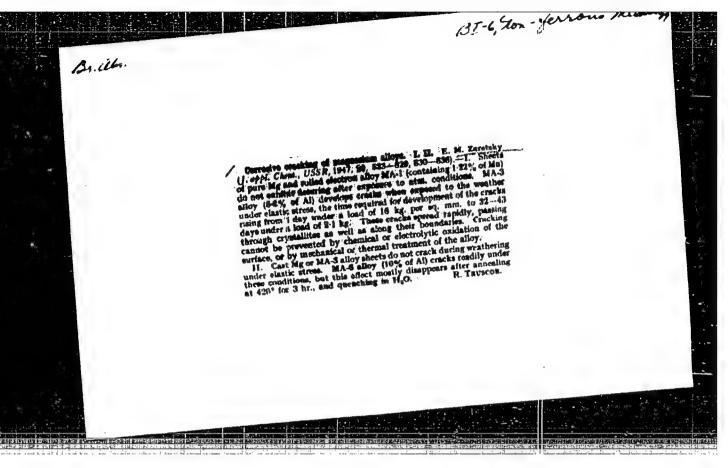
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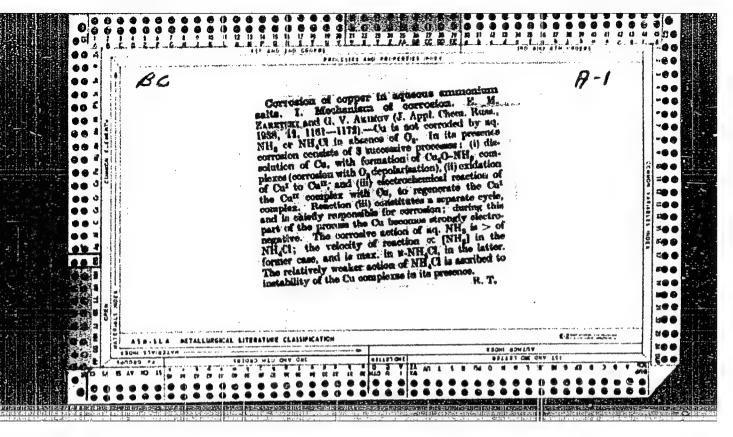
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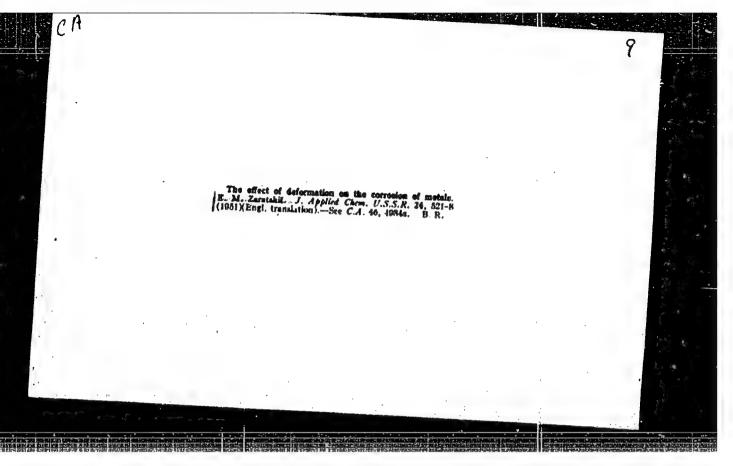


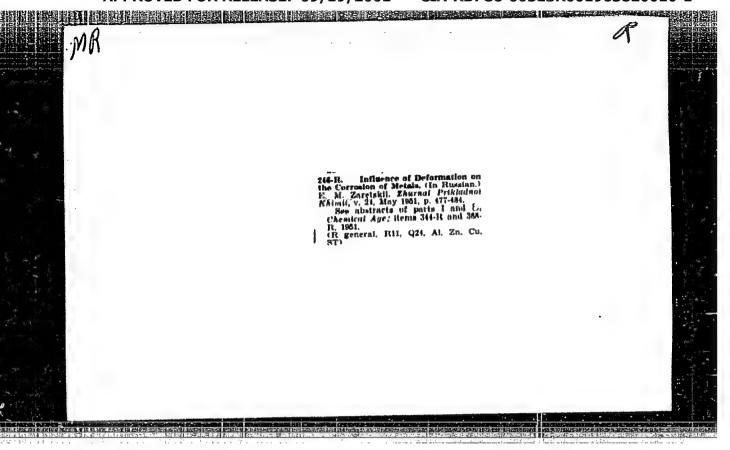


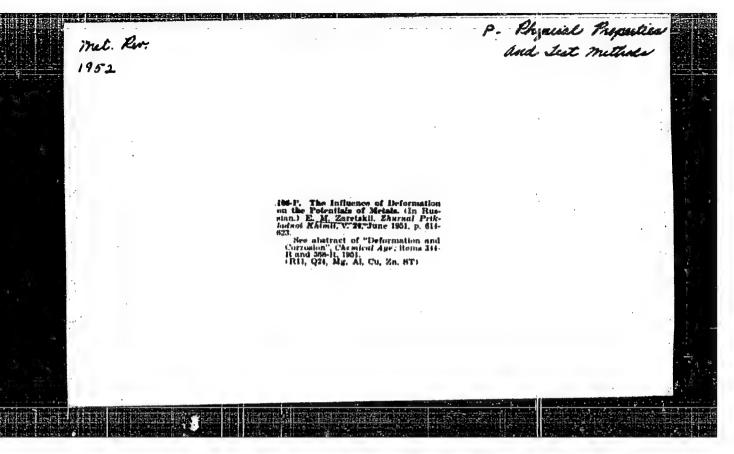


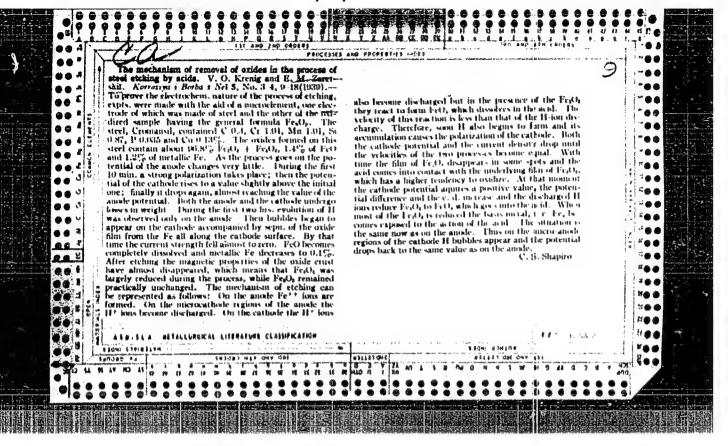


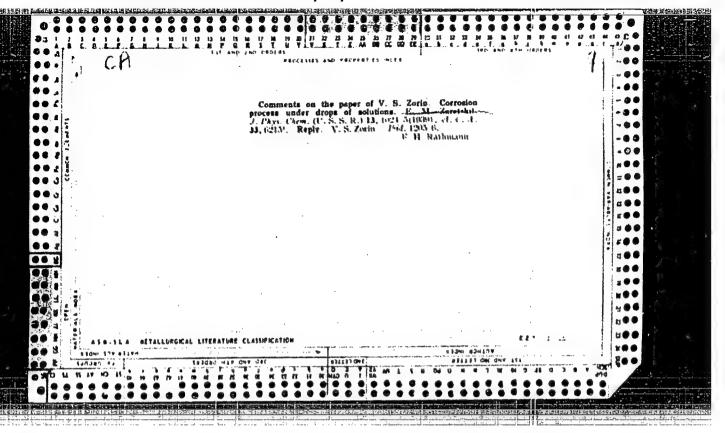


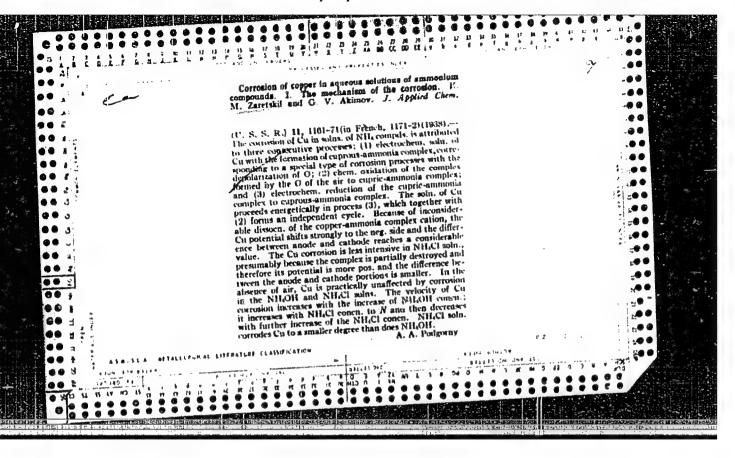


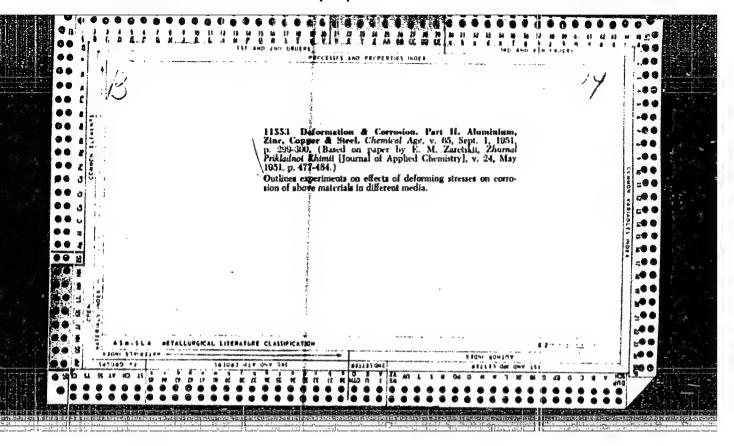


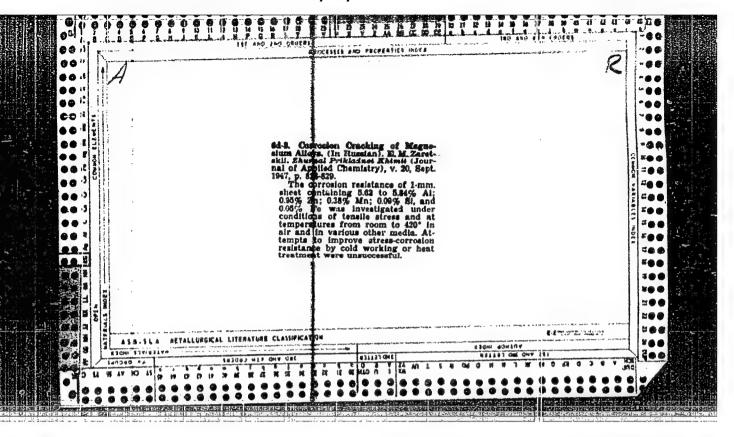


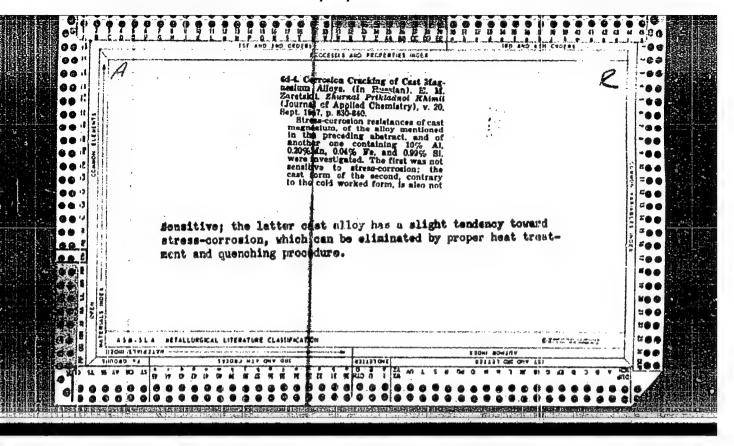


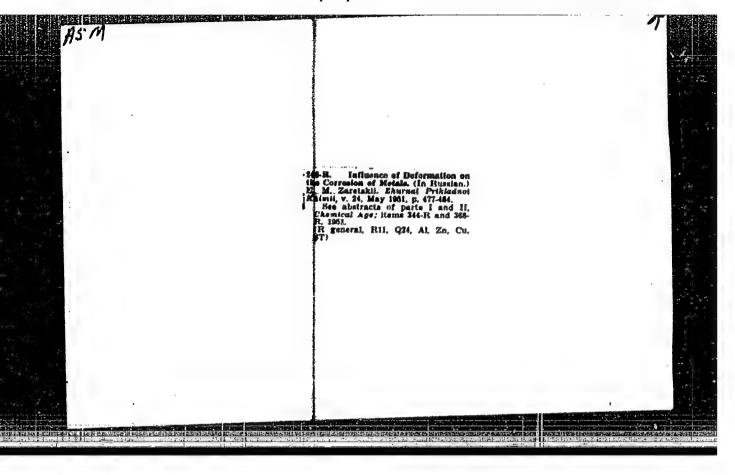


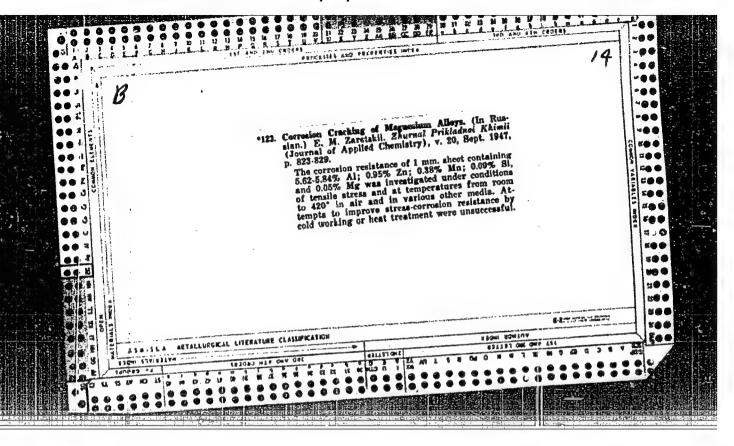


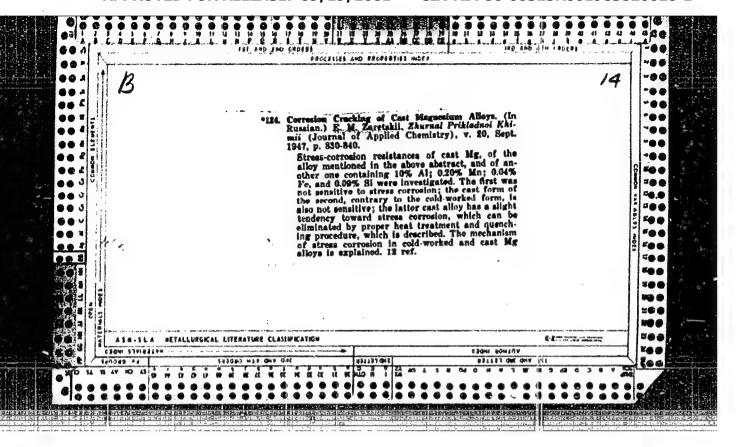


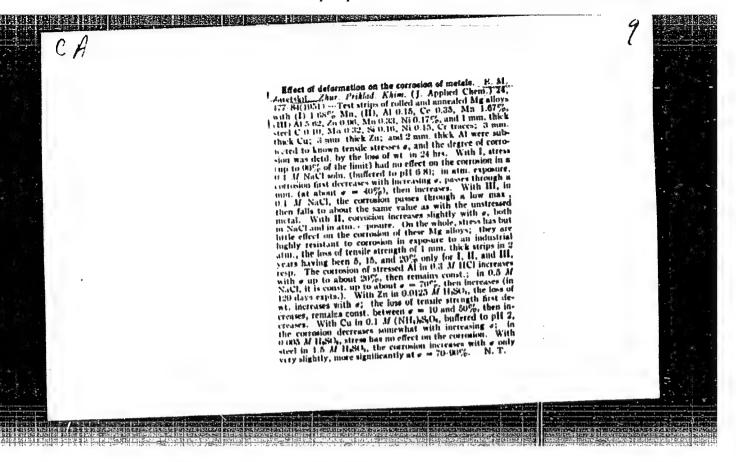


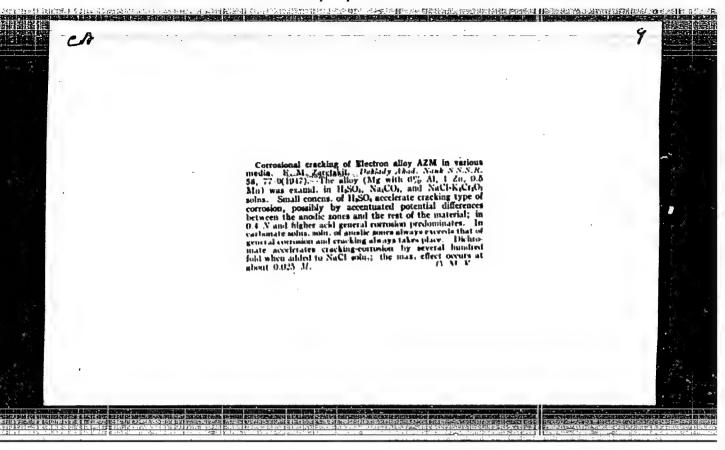


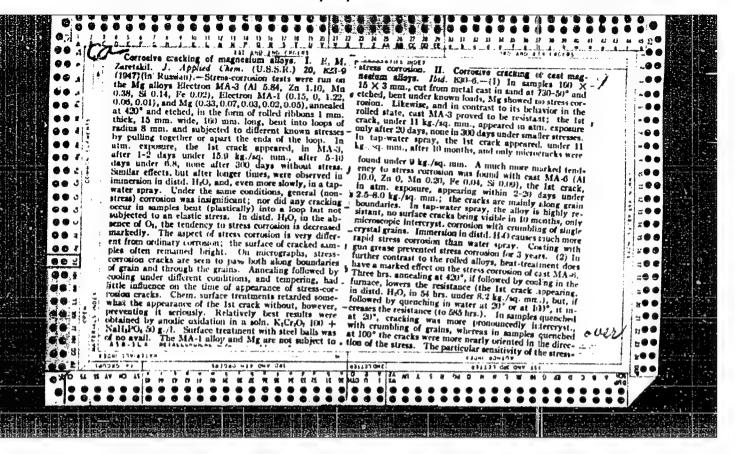


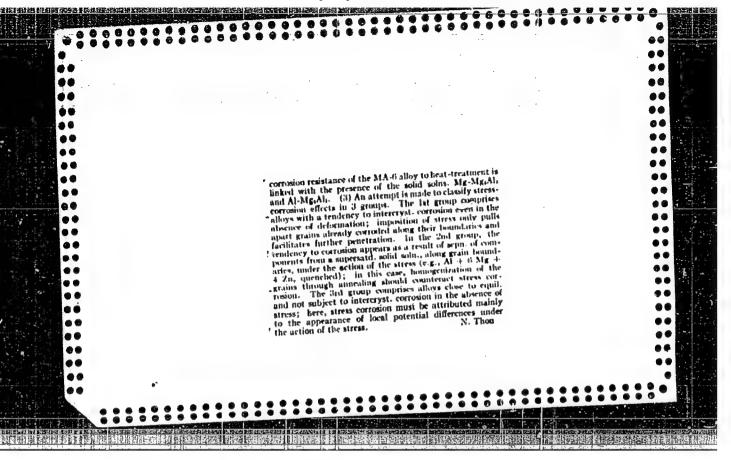


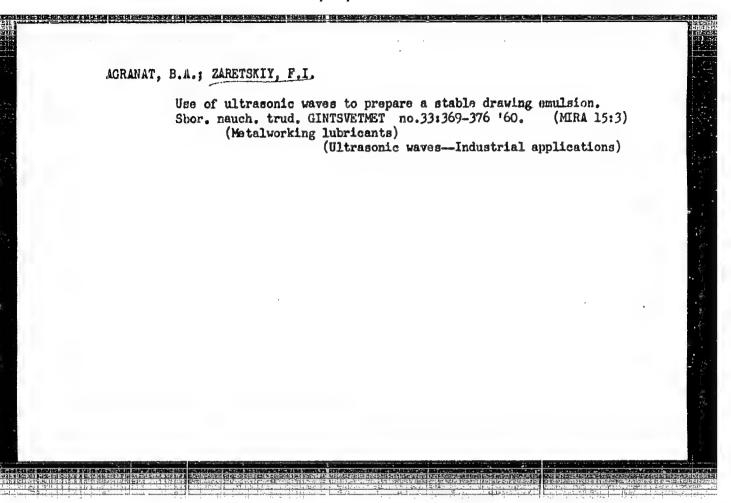












s/137/62/000/c04/068/201 A052/A101

AUTHORS:

Agranat, B. A., Zaretskiy, F. I.

TITLE:

Application of ultrasound to the production of a stable drawing

emulsion

PERIODICAL: Referativnyy zhurnal, Metallurgiya, no. 4, 1962, 34, abstract 4D197

("Sb. nauchn. tr. In-t tsvet. met. im. M. I. Kalinina", no. 33,

1960. 369-376)

The investigation of the emulsification process acted upon by a strong ultrasonic field was carried out on emulsion of the "Moskabel'" plant. TEXT: For drawing Cu- and Al-wire the plant used to prepare emulsions by boiling a scap solution and mechanical mixing with oil and water. These emulsions would usually stratify in 5 - 10 days and proved unsuitable for further use. A series of experiments carried out to produce a more stable emulsion by means of ultrasound has shown that the disperisty of the ultrasound treated emulsion changes inconsiderably, and it preserves practically its stability during several months. A description and diagram of the ultrasonic unit for producing emulsion K. Ursova are given.

[Abstracter's note: Complete translation]

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APPROVED FOR RELEASE: 09/19/2001 CIA-RDP86-00513R001963820010-1"

KOVALENKO, D.N.; ZARETSKIY, G.H.

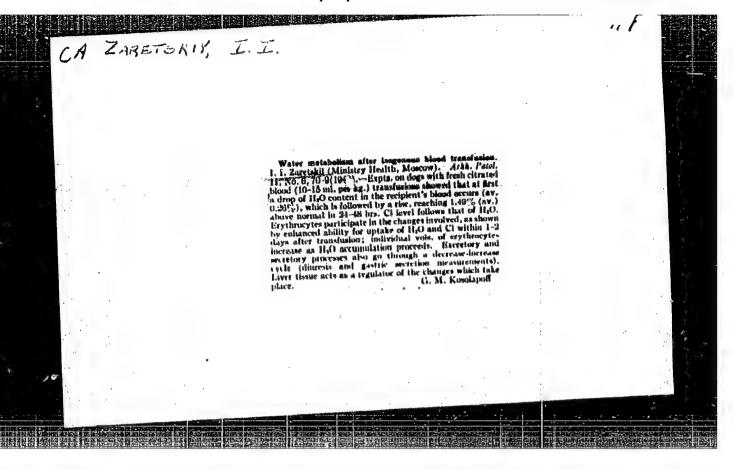
Bentonites in the Lopushnianskoye deposit. Bent.glimy Ukr.
no.3:3-13 '59. (MIRA 12:12)

1. Institut geologicheskikh nauk AN USSR i Ukrainskoye
otdeleniye Veseoyuznogo tresta stroitel'no-tekhnicheskikh
izyskaniy. (Dniester Valley-Bentonite)

SEURKOVICH, S.V.; ZARRTSKIY, I.I.

Effect of immunotherapy on kidney function in burns. Khirurgiia 35 no.7:16-21 J1 '59. (MIRA 12:12)

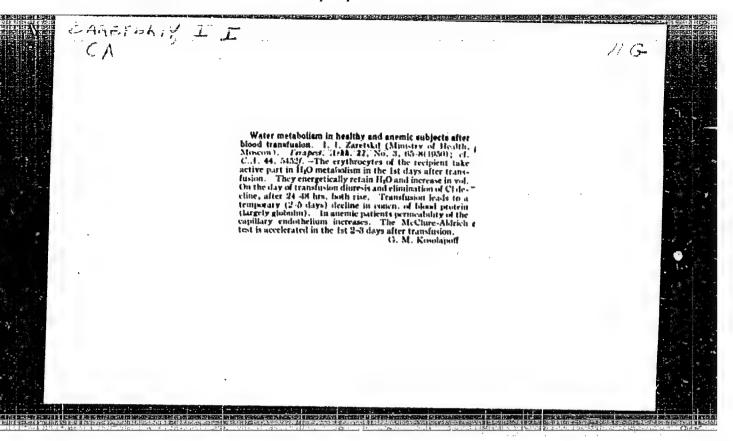
1. Iz patofiziologicheskoy laboratorii (zav. - chlen-korrespondent ANN SSSR prof. N.A. Fedorov) TSentral'nogo instituta genatologii i perelivaniya krovi (dir. - deystvitel'nyy chlen AMN SSSR prof. A.A. Bagdasarov) Ministerstva zdravokhraneniya SSSR. (BURNS, experimental) (KIDNEYS, physiology) (IMMUNE SEHUMS, pharmacology)



ZARETSKIY, I. I.

35481. Izmenehiye Khimicheskogo i morfologicheskogo sostava Krovi vosle
perelivaniya izogennoy kvovi. Vracheb. delo, 1949, No. 11, stb, 985-88.

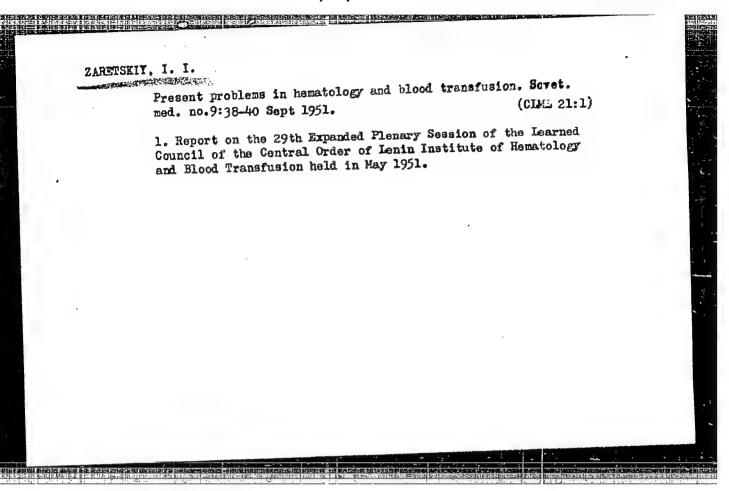
Letopis' Zhurnal'nykh Statey, Vol. 48, Moskva, 1949



ZARETSKIY, I. Candidate of Medical Science
"Problems in Connection with Elood Transfusion"

Meditsinskiy Rabotnik, No. 23, 1950

-W-11691, 5 Jul 1950



ZARETSKIY, I.I., kandidat meditsinskikh nauk (Moscow).

Certain results and proceects in the development of heratology and blood transfusion. Sov.med. 17 no.8:44-47 Ag '53.

(Blood-Transfusion)

1953 Session of the Scientific Council, Central Order of Lenin Institute of Hematology and Blood Transfusion

ZAHRTSKIY, I. I. (Moscow); MIKHAYLOVA, I. A. (Moscow); ROZANOVA, W. S. (Hoscow)

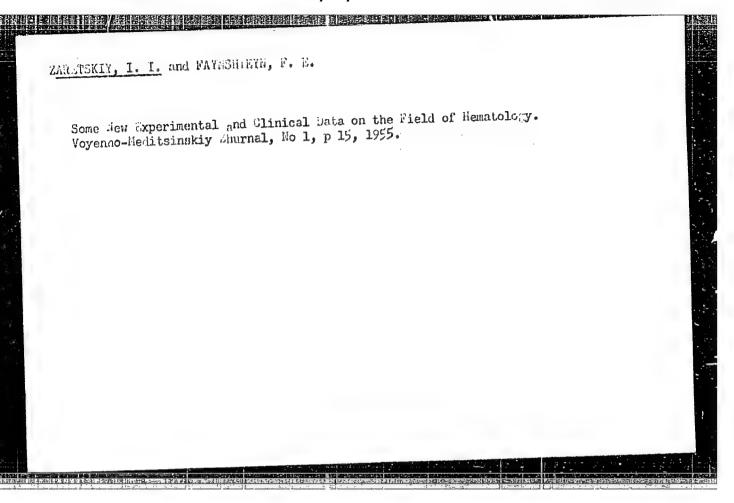
Functional state of the kidneys following transfusion of compatible blood. Arkh.pat. 16 no.2:26-31 Ap-Je '54. (MIRA 7:5)

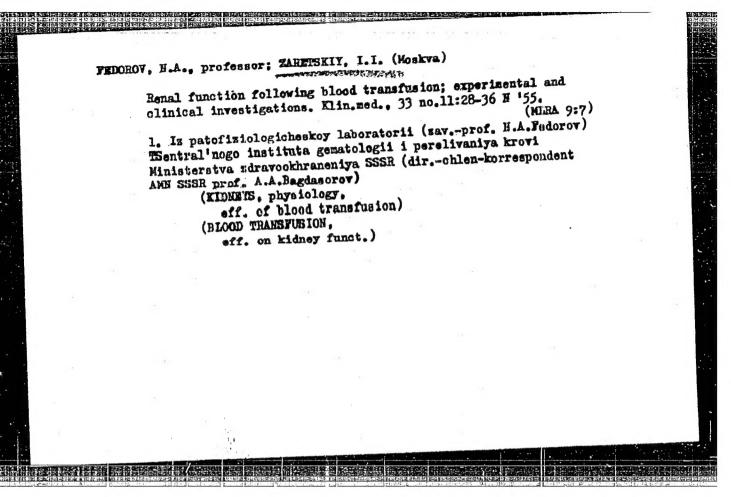
1. Iz patofiziologicheskoy laboratorii (zav. prof. N.A.Fedorov)
TSentral'nogo instituta gematologii i perelivaniya krovi, Ministerstwa zdravockhraneniya SSSR (dir. chlen-korrespondent AMM SSSR prof. A.A.Bagdasarov).

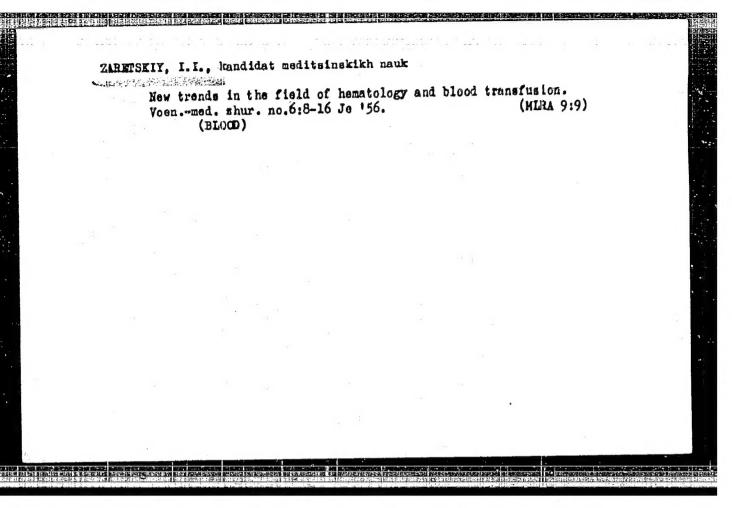
(BLOOD TRANSFUSION. (KIDNAYS, physiology.
*eff. on kidney funct.)

*eff. of blood transfusion)

APPROVED FOR RELEASE: 09/19/2001 CIA-RDP86-00513R001963820010-1"







APPROVED FOR RELEASE: 09/19/2001 CIA-RDP86-00513R001963820010-1"

USSR / Human and Animal Physiology. Excretion.

T

Abs Jour: Ref Zhur-Biol., No 22, 1958, 101947.

Abstract: activity of the tubules according to dioarast or phenol red were determined. The investigations were conducted on the day of transfusion after the emergence of the animals from the state of shock, next day, and later every two days until return to normal. In heterotransfusion and in isotransfusion, a biphase change of the kidney function occurred, sharper in the first case. During the first phase (1-2 days) after blood infusion, the excretory and concentrating powers of the kidneys decreased, glomerulus filtration and the kidneys blood circulation decreased, the filtration fraction and permeability of the glomerular membrane increased. This was confirmed by histologic data (edema of the medular layer of the kidneys, edema

Card 2/3

of the glomerular capsule, constriction of the

VINOKUROVA, G.P.; ZARKTSKIY, I.I.; MIKHAYLOVA, I.A.

The effect of blood transfusion, blood components and plasma substitutes on kidney functions. Problegemat. 1 perel.krovi 1 no.2; 48-52 Mr-Ap '56. (MIRA 10:1)

1. Iz TSentrel'nogo ordena Lenina instituta gematologii i perelivaniya krovi (dir. - chlen-korrespondent AMN SSSR prof. A.A.Badasarov)

Ministerstva săravookhraneniya SSSR.

(KIDDIEYS, physiol.

funct., eff. of blood transfusion, blood components and plasma substitutes)

(BLOUD TRANSFUSION

eff. on kidney funct.)

(PIABMA SUBSTITUTES, eff.
 on kidney funct.)